

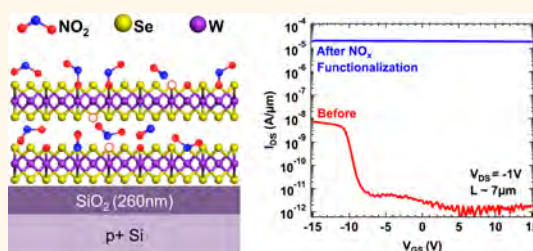
# Air Stable p-Doping of WSe<sub>2</sub> by Covalent Functionalization

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**ABSTRACT** Covalent functionalization of transition metal dichalcogenides (TMDCs) is investigated for air-stable chemical doping. Specifically, p-doping of WSe<sub>2</sub> via NO<sub>x</sub> chemisorption at 150 °C is explored, with the hole concentration tuned by reaction time. Synchrotron based soft X-ray absorption spectroscopy (XAS) and X-ray photoelectron spectroscopy (XPS) depict the formation of various WSe<sub>2-x-y</sub>O<sub>x</sub>N<sub>y</sub> species both on the surface and interface between layers upon chemisorption reaction. *Ab initio* simulations corroborate our spectroscopy results in identifying the energetically favorable complexes, and predicting WSe<sub>2</sub>:NO at

the Se vacancy sites as the predominant dopant species. A maximum hole concentration of  $\sim 10^{19}$  cm<sup>-3</sup> is obtained from XPS and electrical measurements, which is found to be independent of WSe<sub>2</sub> thickness. This degenerate doping level facilitates 5 orders of magnitude reduction in contact resistance between Pd, a common p-type contact metal, and WSe<sub>2</sub>. More generally, the work presents a platform for manipulating the electrical properties and band structure of TMDCs using covalent functionalization.



**KEYWORDS:** layered materials · covalent binding · NO<sub>2</sub> · chemisorption · doping

Transition metal dichalcogenides (TMDCs) have garnered much attention in the recent years due to their unique material properties and potential applications in various electronic and optical devices.<sup>1,2</sup> Possessing a layered structure with a monolayer thickness of down to 7 Å, TMDCs represent a viable electronic material system for future devices, especially as size scaling continues to the sub-10 nm regime. To facilitate TMDC integration into future electronic devices, doping becomes a vital stepping-stone. Indeed, many exploratory works have already reported both n- and p-doping utilizing surface charge transfer mechanisms aided by molecular physisorption.<sup>3–5</sup> For instance, WSe<sub>2</sub> has been p-doped using physisorbed NO<sub>2</sub>; while this doping scheme has achieved superb p-type field-effect transistors (FETs), the doping is unstable and decays almost completely within an hour.<sup>4</sup> The instability is due to the physisorption process, which relies on weak van der Waals (vdW) interactions incapable of sustaining stable electron

transfers at ambient conditions. Moving forward, an *air stable* doping scheme is needed as a fundamental enabler toward the integration of TMDCs into robust device fabrications.

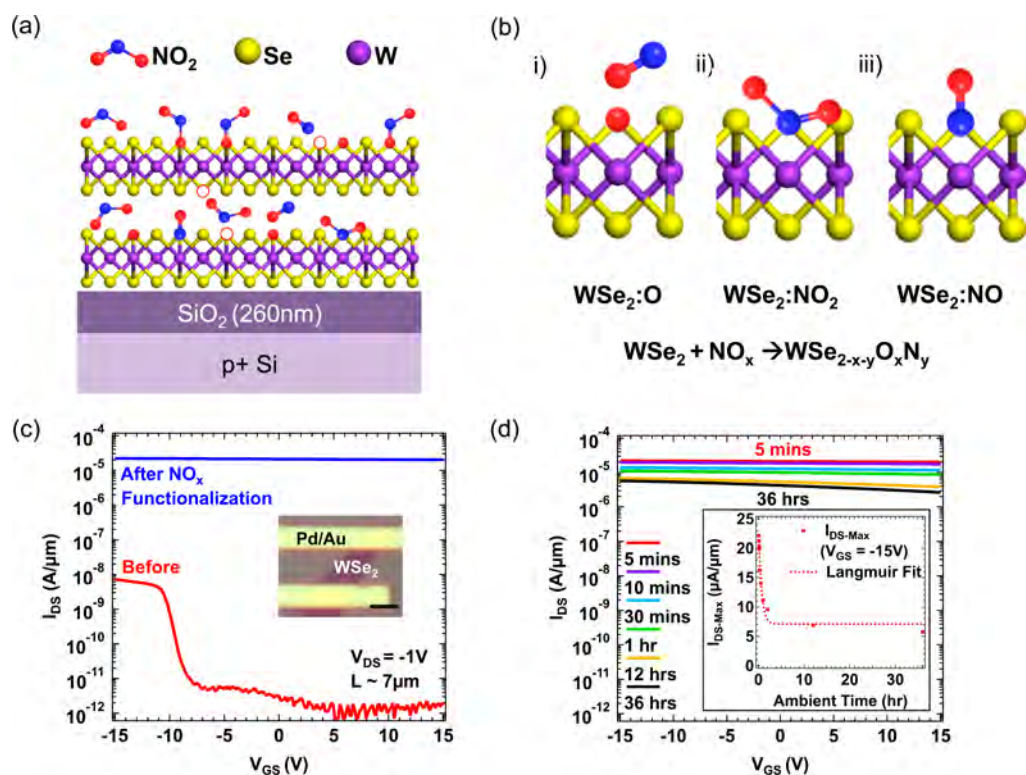
In this work, we demonstrate air stable p-doping, up to the degenerate limit, by *chemisorption* of NO<sub>2</sub> on WSe<sub>2</sub> performed at a sample temperature of 150 °C. Synchrotron-based soft X-ray absorption spectroscopy (XAS) and X-ray photoelectron spectroscopy (XPS) depict formation of various WSe<sub>2-x-y</sub>O<sub>x</sub>N<sub>y</sub> species upon NO<sub>x</sub> chemisorption on WSe<sub>2</sub>. Formation of these species and their effects on the electronic properties of WSe<sub>2</sub> are also examined by *ab initio* simulations, which confirm the experimental observations. Notably, the chemisorption process retains WSe<sub>2</sub>'s long-range crystal structure and optical properties. This doping process leads to a 5 orders of magnitude reduction in contact resistance ( $R_c$ ) between WSe<sub>2</sub> and Pd metal, and results in a degenerate doping concentration of  $1.6 \times 10^{19}$  cm<sup>-3</sup>. From this study, we exhibit a

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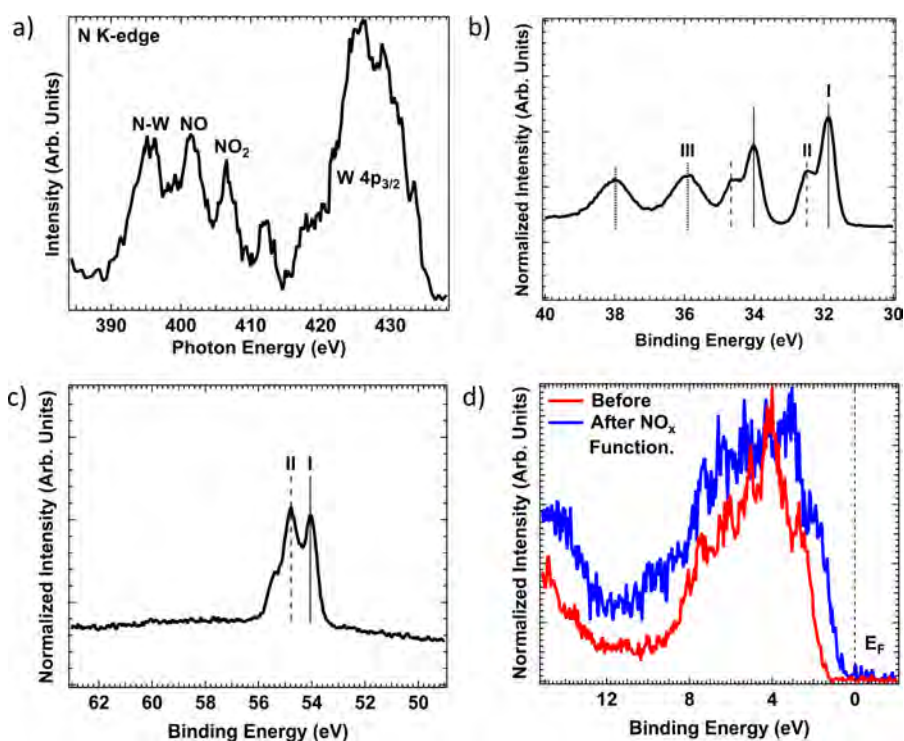
**Figure 1.** (a) Schematic of NO<sub>x</sub> chemisorption process at the WSe<sub>2</sub> bulk and surface. (b) Proposed specific NO<sub>x</sub> chemisorption at the selenium vacancy sites can lead to three distinct configurations: (i) WSe<sub>2</sub>:O, (ii) WSe<sub>2</sub>:NO<sub>2</sub>, (iii) WSe<sub>2</sub>:NO. (c)  $I_{DS}$ – $V_{GS}$  of before and after NO<sub>2</sub> treated devices. Inset: Optical microscopy of fabricated device using Pd/Au contacts; scale bar is 2 μm. (d) Initial current drop after doping is attributed to desorption of weakly physisorbed NO<sub>2</sub>. Inset: Drain current change is fitted to a time-dependent Langmuir isotherm desorption model vs air exposure time. Dash line represents fitted curve.

route toward an air stable p-doping method on WSe<sub>2</sub> that can be efficiently utilized for various device schemes in the future.

## RESULTS AND DISCUSSION

A back-gated WSe<sub>2</sub> FET is fabricated *via* standard electron beam lithography or photolithography and exposed to NO<sub>2</sub> at 150 °C for a set reaction time of 4–12 h. Given the strong oxidizing nature of NO<sub>2</sub>, we predict an induced NO<sub>x</sub> chemisorption process on WSe<sub>2</sub> due to the temperature and treatment time utilized. A defect-oriented model is proposed where NO<sub>2</sub> chemisorbs on WSe<sub>2</sub> surface and bulk defect sites (*e.g.*, selenium vacancies), forming stable electron withdrawing WSe<sub>2-x-y</sub>O<sub>x</sub>N<sub>y</sub> species that lead to p-doping (Figure 1a,b). At the temperatures used here, NO<sub>2</sub> gas is partly decomposed to form O and NO,<sup>6,7</sup> making various covalent functionalization schemes possible. Figure 1b depicts three most likely scenarios. Figure 1b (i) shows direct W oxidation resulting from the O of NO<sub>2</sub> occupying a Se vacancy followed by NO<sub>2</sub> thermal disassociation and desorption of NO. Figure 1b (ii) shows the alternative NO<sub>2</sub> adsorption configuration with N directly bonding to W at the presence of a Se vacancy. Figure 1b (iii) illustrates N of NO covalently bonded with W, where NO is created through NO<sub>2</sub> disassociation.

Transfer characteristics of doped devices are measured before and after NO<sub>x</sub> functionalization (Figure 1c). Here the thickness of WSe<sub>2</sub> flake is ~5 nm. Before NO<sub>2</sub> exposure, we observe a clear p-FET characteristic with an on/off ratio of ~1000 and a low  $I_{on}$  of ~10<sup>-8</sup> A/μm due to severe contact resistance resulting from the Schottky barrier present between WSe<sub>2</sub> and Pd.<sup>4</sup> After NO<sub>x</sub> chemisorption, we observe a drastic increase of  $I_{on}$  by about 1000×, and gate's inability to control the channel conduction. Gate control loss is due to hole doping by NO<sub>x</sub> chemisorption that moves the Fermi level ( $E_F$ ) to close proximity of the valence band edge ( $E_V$ ). The excessive states present in the vicinity of  $E_F$  then render gate modulation of  $E_F$  ineffective, reflecting a degenerate p-doping situation. The doping stability is studied by examining the changing transfer characteristics of the doped sample upon air exposure. As seen in Figures 1d and Supporting Information Figure S1, the device undergoes a gradual current drop for the first 2–3 h, and eventually stabilizes at a set current level. This initial current drop is attributed to desorption of physisorbed NO<sub>x</sub> species leaving the WSe<sub>2</sub>. As expected, our doping mechanism is a combination of NO<sub>x</sub> physisorption and chemisorption on WSe<sub>2</sub> where *only* chemisorbed NO<sub>x</sub> contributes to stable doping. The falling current *versus* time is fit to Langmuir isotherm assuming current being directly related to the physisorbed NO<sub>x</sub> *via*



**Figure 2.** (a) Total electron yield of Soft X-ray Absorption Spectroscopy (XAS)  $\text{NO}_2$  doped samples. Identifiable peaks N K-edge indicates N–W, NO, and  $\text{NO}_2$  species in host after  $\text{NO}_x$  chemisorption. (b) Tungsten core level via X-ray Photoelectron Spectroscopy (XPS) showing the formation of various  $\text{WSe}_{2-x-y}\text{O}_x\text{N}_y$  species. Labels I, II, and III indicate W doublets of different oxidation states. (c) Se 3d core level after  $\text{NO}_x$  chemisorption. Labels I and II indicate two different species of different Se oxidation states. (d) Valence band edge of  $\text{WSe}_2$  before and after  $\text{NO}_x$ ; a definitive  $\sim 0.8$  eV shift of the valence band edge is shown with respect to the Fermi level.

$I(t) = \{I_{\text{initial}} + [\psi/(1 + \psi)][1 - e^{-(t/\tau)}]\}$  where  $t$  is the air exposure time,  $I_{\text{initial}}$  is the initial current level right after chemisorption,  $\psi$  is a dimensionless parameter characterizing surface coverage, and  $\tau$  is the desorption time constant. We extract the desorption time constant  $\tau \sim 40$  min from  $\text{NO}_x$  physisorption, comparable to other physisorption-based gas sensors (such as ion-sensitive FET sensors) operating at room temperature.<sup>8</sup> Notably, the doping effect is irreversible even after thermal annealing of the sample at  $300^\circ\text{C}$  in Ar, which further supports our proposed scheme of  $\text{NO}_x$  chemisorption leading to stable doping.

Given the stability of chemisorption, in-depth material characterization techniques in ultra-high vacuum (UHV) environments can be used to investigate the nature of chemisorbed species, in contrast with that of the unstable surface physisorbed molecules.<sup>4,5</sup> First, the nitrogen K absorption edge was examined using XAS. A collection of both total electron yield (penetration depth of  $\sim 10$  nm) and total fluorescence yield (penetration depth of  $\sim 100$  nm) show similar emission trend, signifying various nitrogen containing species in both depth regimes. The results suggest that the chemisorption is occurring at both the surface and interface between the layers of  $\text{WSe}_2$ . As a function of absorbed photon energy for N 1s electrons, a larger energy translates to a higher N oxidation state, as it becomes harder to remove electrons due to stronger

electrostatic interactions. We observe 3 peaks at  $\sim 396$ ,  $\sim 401$ , and  $\sim 406$  eV which we assign to N–W (from  $\text{W-NO}_x$  species), NO, and  $\text{NO}_2$  species, respectively, based on literature values.<sup>9–14</sup> Therefore, XAS results substantiate all three predicted configurations shown in Figure 1b.

Core W levels were investigated via XPS. According to Figure 2b, six visible peaks were observed after doping compared to only two visible peaks before  $\text{NO}_x$  chemisorption. The six peaks correspond to three doublets (each doublet possessing a unique W  $4f_{5/2}$  and W  $4f_{7/2}$  peak value), each doublet indicating a different W containing species with W of different oxidation states. Doublet III at 35.81 eV is assigned to  $\text{WO}_x$  with at an almost  $6^+$  oxidation state.<sup>15,16</sup> Doublets I and II (31.8 and 32.5 eV for W  $4f_{7/2}$  values, respectively) are assigned to different stoichiometric ratio of  $\text{WSe}_{2-x-y}\text{O}_x\text{N}_y$  species (also observable at the Se 3d level in Figure 2c) and corresponds to W oxidation state between 0 and  $4^+$  (ref 17). Note that the N stoichiometric ratio  $y$  is less than 0.003, corresponding to an atomic concentration of 0.1% and falling below the XPS detection limit. XAS, with a higher sensitivity due to Synchrotron excitation source, however, clearly identifies the N present and underlines the importance of using both XAS and XPS to probe the complete composition of the chemisorbed species present on  $\text{WSe}_2$ .

The valence band edge of  $\text{WSe}_2$  was also examined by XPS. We observe a 0.8 eV shift of  $E_F$  (Figure 2d) upon

$\text{NO}_x$  chemisorption with  $E_V - E_F \sim 50$  meV extracted via linear extrapolation of the valence band tail. This large  $E_F$  shift is indicative of p-doping of  $\text{WSe}_2$  and suggests a degenerate doping level ( $E_V - E_F < 3kT$ ) by  $\text{NO}_x$  functionalization. From  $E_V - E_F$ , we extract a hole concentration of  $p \sim 1.3 \times 10^{19} \text{ cm}^{-3}$  by using the Joyce Dixon approximation<sup>18</sup> using an effective density of states<sup>19</sup> of  $N_V$  at  $2.54 \times 10^{19} \text{ cm}^{-3}$ .

Subsequently, *ab initio* simulation (see Methods for details) was used to examine the thermodynamic stability of  $\text{WSe}_{2-x-y}\text{O}_x\text{N}_y$  species after  $\text{NO}_x$  functionalization and the changing band structure of  $\text{WSe}_2$  as a result. From density functional theory (DFT) calculations, we observe energetically favorable adsorption of O, NO and  $\text{NO}_2$  at Se vacancies (Table 1), which is consistent with the experimental observations. Notably, on the basis of DFT calculations, adsorption was found to be unfavorable without the presence of Se vacancies.

Out of all adsorbates, only NO adsorption is predicted to contribute to p-doping of  $\text{WSe}_2$  based on DFT simulations. Figure 3 shows the relaxed configuration of chemisorbed NO with N bonding to W, similar to Figure 1b (iii) configuration. By examining the band structure, we find that  $\text{NO}:\text{WSe}_2$  configuration creates two defect bands close to the valence band edge

shown in Figure 3b. The red dash line, overlapping with one of the NO adsorption induced defect bands, represents  $E_F$ . The other two chemisorbed species have minimal effects on the carrier concentration of  $\text{WSe}_2$  based on the DFT simulations (Figure 3c).

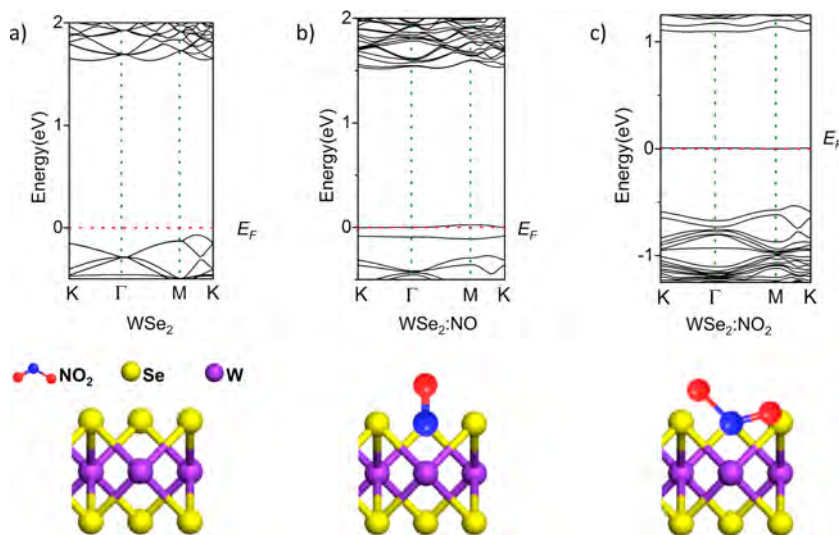
Next we examine the effects of  $\text{NO}_x$  functionalization on the crystal structure and luminescence properties of  $\text{WSe}_2$ . X-ray diffraction (XRD), Raman spectroscopy and photoluminescence (PL) were all performed on  $\text{NO}_x$  functionalized  $\text{WSe}_2$  samples. XRD data in Figure 4a demonstrates that the overall  $\text{WSe}_2$  crystal structure remains intact as the two Miller index peaks 002 and 103 are preserved with similar intensities and no visible  $\theta$  shifts. Figure 4b demonstrates the retention of the primary phonon modes  $E_{2g}$  and  $A_{1g}$  at 249 and 251  $\text{cm}^{-1}$ , respectively, for a trilayer  $\text{WSe}_2$  before and after doping, consistent with the literature.<sup>20,21</sup> PL spectra in Figure 4c shows the two peaks (1.47 and 1.6 eV) of a trilayer  $\text{WSe}_2$  unaffected by  $\text{NO}_x$  chemisorption.<sup>22,23</sup> From our XRD, Raman and PL characterizations, we conclude that doping functionalization does not compromise the unique properties of  $\text{WSe}_2$  at the macroscopic level.

Electrical measurements of  $\text{WSe}_2$  devices before and after  $\text{NO}_x$  chemisorption were performed. Two studies were conducted, the first being a time dependent doping study at varying thicknesses. Figure 5a shows a representative transfer characteristic of 7.5 nm thick  $\text{WSe}_2$  back-gated device as a function of  $\text{NO}_x$  doping time. A progressively heavier p-doping can be observed, reflected by an increasing  $I_{\text{on}}$  and diminishing gate control. The increasing  $I_{\text{on}}$  displays a degree of doping tunability *via* doping time, with  $I_{\text{on}}$  eventually saturating to degenerate behavior. Doping concentration is explicitly extracted across three devices with different flake thicknesses of 0.7, 7.5, and 78 nm

**TABLE 1. Binding Energies of Three Adsorbates ( $\text{NO}_2$ , NO, O) on Se Vacancies of  $\text{WSe}_2$ <sup>a</sup>**

adsorbates on Se vacancies		
$\Delta E_{\text{NO}_2:\text{WSe}_2}$	$\Delta E_{\text{NO}:\text{WSe}_2}$	$\Delta E_{\text{O}:\text{WSe}_2}$
-1.0 eV	-2.8 eV	-4.8 eV

<sup>a</sup> The error for binding energy is within 0.1 eV.



**Figure 3.** *Ab initio* simulation results showing band structure and corresponding simulated model of (a) pristine  $\text{WSe}_2$ , (b)  $\text{WSe}_2$  with NO adsorption in Se vacancy, and (c)  $\text{WSe}_2$  with  $\text{NO}_2$  adsorption in Se vacancy. The red dotted line indicates the Fermi level.



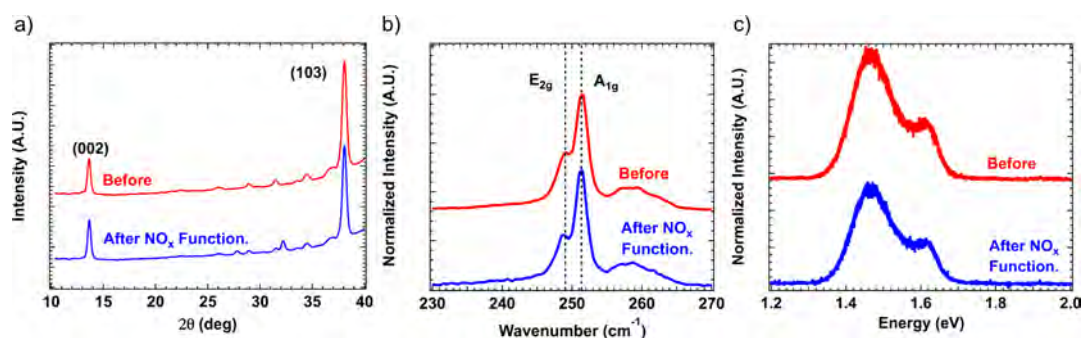


Figure 4. (a) X-ray diffraction (XRD) pattern of WSe<sub>2</sub> before and after WSe<sub>2</sub> functionalization. (b) Raman and (c) photoluminescence (PL) spectra of the same trilayer WSe<sub>2</sub> flake after NO<sub>x</sub> chemisorption.

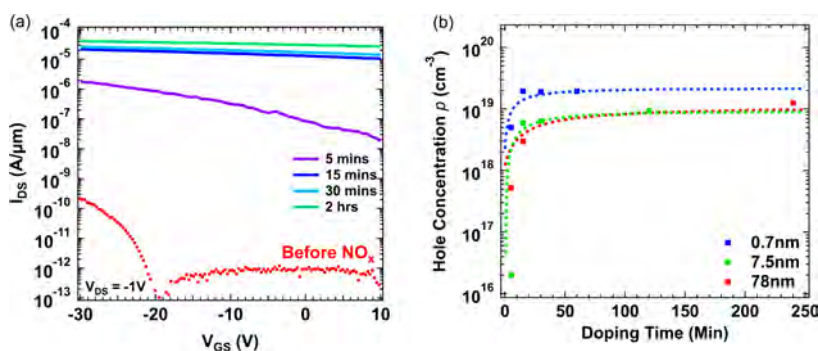


Figure 5. (a) Representative  $I_{DS}-V_{GS}$  for WSe<sub>2</sub> device exposed to NO<sub>2</sub> at different functionalization time. (b) Extracted doping concentration vs doping time for different flake thicknesses. All doping levels saturate at  $\sim 10^{19} \text{ cm}^{-3}$ .

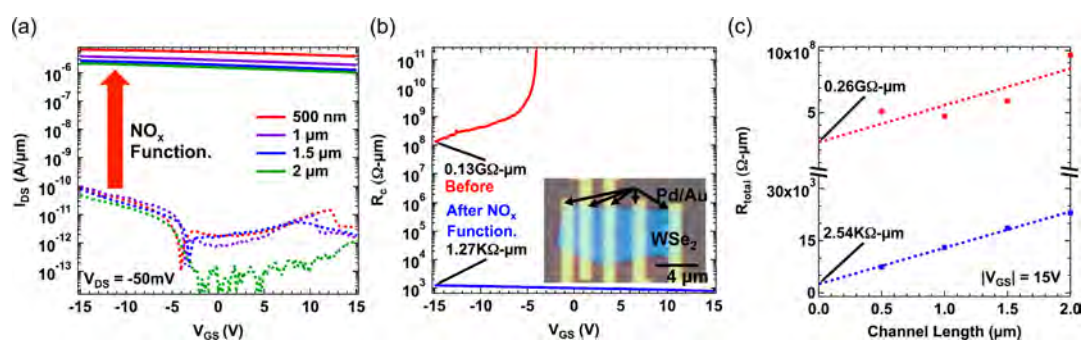


Figure 6. (a)  $I_{DS}-V_{GS}$  of TLM device before and after NO<sub>2</sub> treatment, showing a (b) 5 orders of magnitude decrease in contact resistance. Inset: Optical microscopy of fabricated back-gated device used for TLM. (c) Extracted total resistance change before and after NO<sub>x</sub> functionalization at  $|V_{GS}| = 15 \text{ V}$ .

(monolayer thickness is determined by optical microscopy and the rest determined by AFM in Supporting Information Figure S2), respectively, versus doping time (Figure 5b). Assuming negligible  $R_c$  due to the thinning Schottky barrier at higher doping concentrations and therefore higher tunneling current, we can estimate doping concentration by first finding the 2D sheet charge density  $n_{2D} = (I_{DS}L)/(qWV_{DS}\mu_p)$  at  $V_G = 0$  where  $V_{DS}$  is the drain voltage,  $L$  and  $W$  are the length and width of the channel,  $q$  is the unit electron charge, and  $\mu_p$  is the field effect hole mobility. From  $n_{2D}$ , we find volume doping concentration which saturates to  $\sim 1 \times 10^{19} \text{ cm}^{-3}$  across all three devices. The similar saturated doping level across all three devices demonstrates a uniform volume doping throughout the WSe<sub>2</sub> surface and body, supporting the XAS results discussed

earlier. Notably, the electrically extracted saturation doping concentration is consistent with that extracted from XPS measurements.

Contact resistance,  $R_c$  was extracted before and after NO<sub>x</sub> chemisorption using the transfer line method (TLM).<sup>24</sup> Figure 6a shows the transfer characteristics of TLM devices with channel lengths of 500 nm, 1  $\mu\text{m}$ , 1.5  $\mu\text{m}$ , and 2  $\mu\text{m}$  fabricated on a WSe<sub>2</sub> flake with a thickness of  $\sim 7 \text{ nm}$  (Figure 6b inset). The observed low  $I_{on}$  and ambipolar behavior prior to doping implies a large  $R_c$  due to Schottky barrier height near the WSe<sub>2</sub> midgap. After doping, however, we see a significant increase in  $I_{on}$  indicating a decrease in  $R_c$ . Using  $R_{total} = 2R_{contact} + R_{channel}(L)$ , where  $R_{total}$  and  $R_{channel}$  are the unit-width normalized total resistance and channel resistance,  $R_c$  is found to be 0.13  $\text{G}\Omega/\mu\text{m}$  and 1.27

$k\Omega/\mu\text{m}$  before and after  $\text{NO}_x$  chemisorption, respectively (Figure 6b,c).  $\text{NO}_2$  doping is shown to reduce the  $R_c$  by a remarkable 5 orders of magnitude. Finally, we revisit  $\text{WSe}_2$  doping extraction including  $R_{c1}$  and find again consistent doping level  $p \sim 1.6 \times 10^{19} \text{cm}^{-3}$ .

## CONCLUSION

In summary, we have demonstrated air stable p-doping of  $\text{WSe}_2$  via covalent functionalization by  $\text{NO}_x$  at an elevated temperature of  $150^\circ\text{C}$ . The stability of this covalent bonding scheme enables in-depth characterization via XAS and XPS, confirming  $\text{WSe}_{2-x-y}\text{O}_x\text{N}_y$  species formation due to  $\text{NO}_x$  chemisorption on both the  $\text{WSe}_2$  surface and body. From XPS, a shift of  $\sim 0.8$  eV of  $E_F$  toward  $E_V$  also concludes a degenerate doping situation. In addition, all observed

$\text{WSe}_{2-x-y}\text{O}_x\text{N}_y$  species are shown to be energetically favorable from *ab initio* simulations, with  $\text{NO}:\text{WSe}_2$  identified as the dominant configuration leading to p-doping. Our time dependent doping study also indicates doping tunability, and the TLM study further confirms the extracted doping concentration of  $1.6 \times 10^{19} \text{cm}^{-3}$  while also showing a remarkable  $\sim 5$  orders of magnitude  $R_c$  reduction via  $\text{NO}_x$  functionalization. Controllable semiconductor doping is an important and fundamental enabler in the general field of electronic devices, and degenerate materials can be utilized anywhere from contact materials to junctions in tunneling devices. Therefore, our demonstration of air stable doping through covalent functionalization of  $\text{WSe}_2$  presents another route for further enhancing and manipulating TMDC functionalities and applications.

## METHODS

**Device Fabrication.** The device fabrication starts with mechanical exfoliation of  $\text{WSe}_2$  (Nanosurf) on top of a  $\text{SiO}_2$  (260 nm)/Si substrate. Source/drain contacts are patterned using photolithography, and e-beam evaporation of Pd/Au (10/30 nm) films and resist liftoff in acetone.

**$\text{NO}_2$  Doping.** The fabricated sample is placed inside an evacuated 250 mL Pyrex glass flask with ambient  $\text{N}_2$  flow for gas exchange. Subsequently, a constant  $\text{NO}_2$  gas flow (5000 ppm balanced in  $\text{N}_2$ ) is introduced through the flask and maintained throughout the entire doping procedure, exhausting to an external oil bath. Simultaneously, the flask is placed on a hot plate with the sample temperature maintained at  $150^\circ\text{C}$  and doped for a set time. Finally, the sample is cooled first to  $\sim 50^\circ\text{C}$  before purging the flask with  $\text{N}_2$  and sample removal.

**XAS Characterizations.** X-ray absorption spectroscopy of  $\text{NO}_x$  chemisorption in  $\text{WSe}_2$  characterization was taken at Beamline 8.0.1 at Advanced Light Source, Lawrence Berkeley National Laboratory. N K-edge spectrum was calibrated carefully by measuring h-BN reference sample for energy calibration. The energy resolution at N K-edge region is about 0.1 eV and the vacuum of experiment chamber is lower than  $5 \times 10^{-9}$  Torr. All XAS spectra were recorded in Total Electron Yield (TEY) and Total Fluorescence Yield (TFY) simultaneously.

**XPS Characterizations.** The XPS spectra were acquired using a monochromated Al K $\alpha$  source ( $h\nu = 1486.7$  eV) connected to a UHV cluster system described elsewhere.<sup>25</sup> An analyzer acceptance angle of  $\pm 8^\circ$ , a take-off angle of  $45^\circ$ , and pass energy of 15 eV were used for this study. The binding energy scale is calibrated using the ASTM procedure.<sup>26</sup>

**Optical, Electrical, Raman, PL, and XRD Characterizations.** Optical images are taken with an Olympus BX51 microscope with an attached video camera. All electrical measurements are conducted on a HP 4155C analyzer with a corresponding probe station. Both Raman and PL are conducted using a commercially available HORIBA LabRAM HR800 system using an external laser source with an excitation wavelength of 532 nm. The X-ray diffraction (XRD) was taken on a Bruker AXS D8 Discover GADDS XRD Diffractometer system.

**Ab Initio Simulations.** First-principles calculations were utilized based on DFT<sup>27,28</sup> with plane wave basis sets and Projector Augmented Wave (PAW) pseudopotentials<sup>28,29</sup> implemented in the Vienna *Ab initio* Simulation Package (VASP).<sup>27,30,31</sup> The electronic wave functions were represented by plane wave basis with a cutoff energy of 600 eV. The exchange and correlation interactions are incorporated as a functional of the Generalized Gradient Approximation (GGA).<sup>29,32,33</sup> A  $5 \times 5$  supercell of monolayer  $\text{WSe}_2$  host was used for the simulations. Each model has  $\sim 18 \text{ \AA}$  vacuum to avoid interaction between replica images

as a result of the periodic boundary conditions.  $\text{NO}$ ,  $\text{NO}_2$  and O were separately placed near a singular Se vacancy of  $\text{WSe}_2$  (or pristine  $\text{WSe}_2$ ) and allowed to relax energetically while the host cell size was kept fixed after optimization.

**Conflict of Interest:** The authors declare no competing financial interest.

**Supporting Information Available:** Detailed device transfer characteristics after doping, and AFM of various flake thicknesses used. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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# Supporting Information

## Air Stable $p$ -Doping of WSe<sub>2</sub> by Covalent Functionalization

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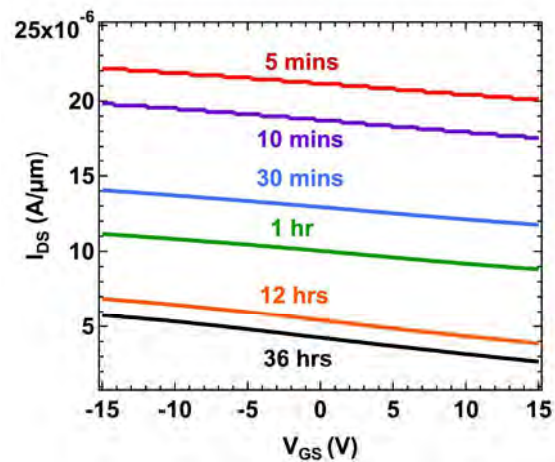
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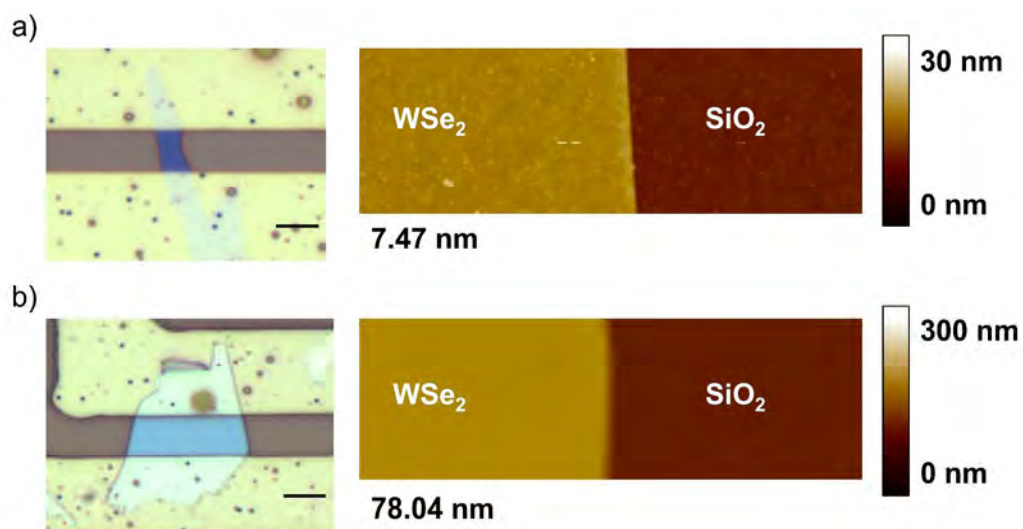
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**Figure S1.** Facsimile of Figure 1(d) represented in linear scale for drain current. A degenerate p-type trend is observed with small on/off ratio and slightly lower drain current at a more positive gate bias.



**Figure S2.** Optical microscopy (left) and atomic force microscopy (AFM, right) images of devices shown in Figure 5. WSe<sub>2</sub> layer thicknesses are (a) ~7.5nm and (b) ~ 78 nm. Scale bar is 3μm.